

Jan 5, 2011

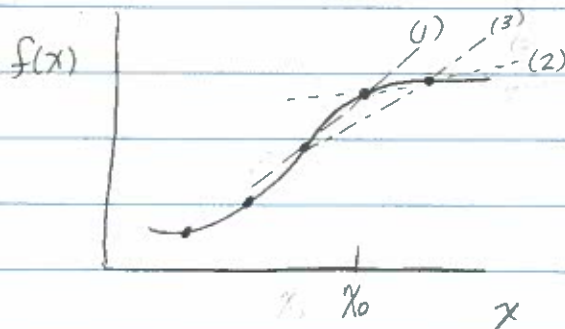
Numerical solutions to wave equations

$$\begin{array}{l}
 (1) \quad \frac{df}{dx}(x_0) = \lim_{\Delta x \rightarrow 0} \frac{f(x_0 + \Delta x) - f(x_0)}{\Delta x} \\
 (2) \quad \frac{df}{dx}(x_0) = \lim_{\Delta x \rightarrow 0} \frac{f(x_0) - f(x_0 - \Delta x)}{\Delta x} \\
 (3) \quad \frac{df}{dx}(x_0) = \lim_{\Delta x \rightarrow 0} \frac{f(x_0 + \Delta x) - f(x_0 - \Delta x)}{2\Delta x}
 \end{array}
 \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} \text{"one"} \\ \text{"sided"} \\ \text{"centered"} \end{array}$$

If $\frac{df}{dx}(x_0)$ is continuous and $\Delta x \rightarrow 0$

all three are identical

But if Δx is finite, (1) - (3) give different approximations to $\frac{df}{dx}$ at x_0 .



points are separated by Δx

When Δx is finite (1) - (3) are called finite differences

e.g.
$$\frac{df}{dx}(x_0) \approx \frac{f(x_0 + \Delta x) - f(x_0)}{\Delta x}$$

What is the accuracy of this approximation?

use Taylor's Series Expansion

$$f(x_0 + \Delta x) = f(x_0) + \Delta x \frac{df}{dx}(x_0) + \frac{\Delta x^2}{2} \frac{d^2f}{dx^2}(x_0) + \dots$$

$$\frac{f(x_0 + \Delta x) - f(x_0)}{\Delta x} - \frac{df}{dx}(x_0) = \frac{\Delta x}{2} \frac{d^2 f}{dx^2}(x_0) + \dots$$

one-sided differences

leading term is of order Δx or $\mathcal{O}(\Delta x)$
 "first-order accurate"

$$\frac{f(x_0 + \Delta x) - f(x_0 - \Delta x)}{2\Delta x} - \frac{df}{dx}(x_0) = \frac{\Delta x^2}{6} \frac{d^3 f}{dx^3}(x_0) + \dots$$

Centered difference are $\mathcal{O}(\Delta x^2)$

Consider the constant wind speed advection eq

$$\frac{\partial \psi}{\partial t} + c \frac{\partial \psi}{\partial x} = 0$$

solution $\psi(x, t) = \psi(x - ct, 0)$ "conveyor belt"

We wish to obtain an approximate solution
 at grid point $(n\Delta t, j\Delta x)$ where n & j are integers

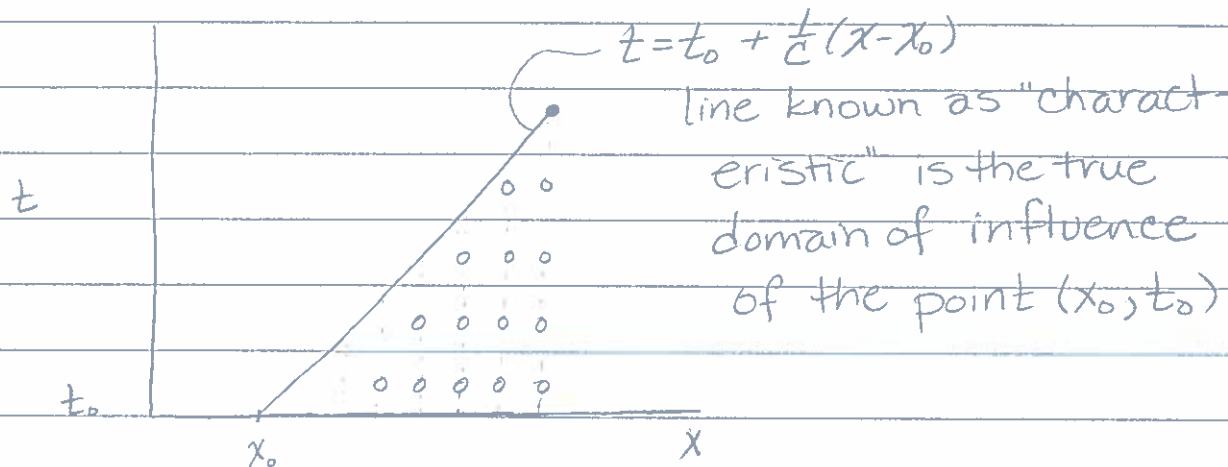
Using $\psi(n\Delta t, j\Delta x) = \psi_j^n$ shorthand

one possible solution is the upstream scheme

$$\frac{\psi_j^{n+1} - \psi_j^n}{\Delta t} + c \frac{\psi_j^n - \psi_{j-1}^n}{\Delta x} = 0$$

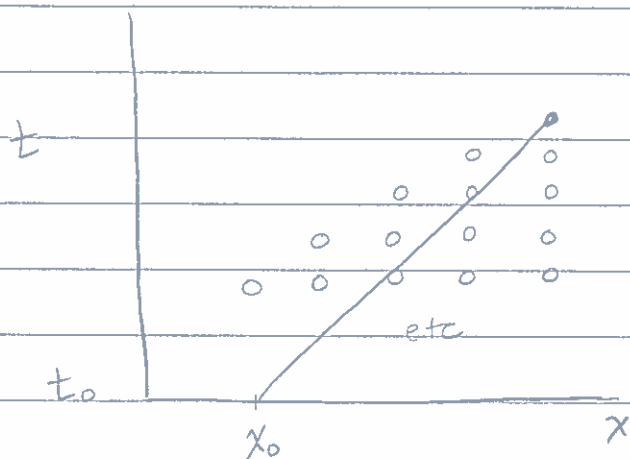
$$\psi_j^{n+1} = \psi_j^n - \frac{c\Delta t}{\Delta x} (\psi_j^n - \psi_{j-1}^n)$$

$$\frac{C\Delta t}{\Delta x} = \text{"Courant" Number}$$



Circles are the domain of dependence of the upstream solution.

The circles must encompass the line if the solution is stable. This is not the case above but is below



What is the key difference?

Δx versus Δt

The stability condition is: $\frac{\Delta t}{\Delta x} < \frac{1}{c}$

The Courant Friedrich Lewy Condition (CFL) for upstream

$$0 \leq C \frac{\Delta t}{\Delta x} \leq 1$$

(the zero is a lower limit because $c, \Delta t, \Delta x$ are each ≥ 0)

A few other solutions are

Forward in time - centered in space

$$\frac{\psi_j^{n+1} - \psi_j^n}{\Delta t} + c \left(\frac{\psi_{j+1}^n - \psi_{j-1}^n}{2\Delta x} \right) = 0$$

Leapfrog in time - centered in space

$$\frac{\psi_j^{n+1} - \psi_j^{n-1}}{\Delta t} + c \left(\frac{\psi_{j+1}^n - \psi_{j-1}^n}{2\Delta x} \right) = 0$$

The different schemes have different behaviors that you will explore in HW1

B.C.'s

Initial Condition

$$\psi_0 = \sin^6 kx$$

B.C. "periodic"

1 \Rightarrow n_x same

2 \rightarrow

centered in space at $n=1$

$$\psi_2^n - \psi_{n_x-1}^n$$