

Combined space/time differencing; consistency

$$L[\psi] = \frac{\partial \psi}{\partial t} + c \frac{\partial \psi}{\partial x} = 0, c > 0 \quad \text{let FDA to } \psi(x, t_n) \text{ be } \phi_j^n$$

Upstream scheme  $\delta_t^F \phi_j^n + c \delta_x^B \phi_j^n = L_a[\phi_j^n] = 0$

$$\begin{aligned} \frac{\partial}{\partial t} &\leftarrow \delta_t^F \\ \frac{\partial}{\partial x} &\leftarrow \delta_x^B \end{aligned}$$

If exact solution is substituted into FDM,

$$\begin{aligned} L_a[\psi^n]_j &= \delta_t^F \psi_j^n + c \delta_x^B \psi_j^n = \Psi_t + \frac{\Delta t}{2} \Psi_{tt} + \dots \\ &\quad + c(\Psi_x + \frac{\Delta x}{2} \Psi_{xx} + \dots) \\ &= \underbrace{\frac{\Delta t}{2} \Psi_{tt} + c \frac{\Delta x}{2} \Psi_{xx}}_{T_a[\Psi]}, \text{ local truncation error.} \end{aligned}$$

If truncation error is  $O(\Delta x^p, \Delta t^q)$  method is  $p$ 'th order accurate in  $t$  and  $q$ 'th order in  $x$ . If  $T_a[\Psi] \rightarrow 0$  as  $\Delta x, \Delta t \rightarrow 0$  the scheme is called consistent.

(D2.2) A FDM is convergent to  $p$ 'th order in  $t$ ,  $q$ 'th order in  $x$  if

$$\|\psi(x, t) - \phi_j^n\| = O(\Delta t^p) + O(\Delta x^q) \quad \text{as } \Delta t, \Delta x \rightarrow 0$$

for all  $x, t_i \leq T$ , where  $T$  is some positive time in  
independent of  $\Delta x, \Delta t$  but, where  $\|\phi\|_p = \left( \sum_{j=1}^N |\phi_j|^p \Delta x \right)^{1/p}$   
( $p=2 \rightarrow$  euclidean norm)  
( $p=\infty \rightarrow$  max norm)  
(Recall properties of norms).

A FDM is stable if

$$\|\phi^n\| \leq C_T \|\phi^0\| \quad \text{for } n \Delta t \leq T, \text{ a fixed positive}$$

Exponentially stable if  $C_T \leq Ae^{BT}$ , strictly stable if  $C_T \leq 1$ .  
time, and all sufficiently small  $\Delta x, \Delta t$ .

Lax Equivalence Thm (Lax & Richtmeyer 1956)

If a FDM is linear, stable and accurate of  $O(\Delta t^p, \Delta x^q)$  it is convergent of  $O(\Delta t^p, \Delta x^q)$

Testing for stability of a FD scheme(i) Energy method

The PDE being approximated often has some form of 'energy' conservation. A well-designed numerical scheme will often preserve this conservation law, which may ensure stability.

E.g.

$$\partial_t \psi_t + c \psi_x = 0, c > 0 \quad \psi(0) = \psi(1)$$

$$\Delta x = \frac{1}{N}$$

$$\text{FDA: } \psi(x_j, t_n) \approx \phi_j^n$$

$$\delta_t^F \phi_j^n + c \delta_x \phi_j^n = 0 \quad , \quad \phi_0^n = 0, \phi_N^n = 0.$$

$$\Rightarrow \phi_j^{n+1} = \mu \phi_{j-1}^n + (1-\mu) \phi_j^n, \mu = \frac{c \Delta t}{\Delta x} = \text{Courant number}$$

Cons. law for PDE:

$$\frac{\partial}{\partial t} \int_0^1 \psi^2 dx = \int_0^1 2\psi \psi_t dx = -2c \int_0^1 \psi \psi_x dx = 0.$$

⇒ Try FDA analogue:

$$\begin{aligned} \frac{\sum_{j=1}^N (\phi_j^{n+1})^2 \Delta x - \sum_{j=1}^N (\phi_j^n)^2 \Delta x}{\Delta t} &= \frac{\Delta x}{\Delta t} \sum_{j=1}^n \left[ \left\{ \phi_j^n - \mu(\phi_j^n - \phi_{j-1}^n) \right\}^2 - \phi_j^n \right] \\ &= \frac{\Delta x}{\Delta t} \sum_{j=1}^n \left[ -2\mu \phi_j^n (\phi_j^n - \phi_{j-1}^n) + \mu^2 (\phi_j^n - \phi_{j-1}^n)^2 \right] \\ &= \frac{\Delta x}{\Delta t} \sum_{j=1}^n \left[ -\mu (\phi_j^{n+2} - \phi_j^{n-2}) - \mu (\phi_j^n - \phi_{j-1}^n)^2 + \mu^2 (\phi_j^n - \phi_{j-1}^n)^2 \right] \\ \frac{\Delta x}{\Delta t} \left\{ \| \phi^{n+1} \|_2^2 - \| \phi^n \|_2^2 \right\} &= -\frac{\Delta x}{\Delta t} \cdot \mu (1-\mu) \sum_{j=1}^n (\phi_j^n - \phi_{j-1}^n)^2 \leq 0 \text{ if } 0 \leq \mu \leq 1 \end{aligned}$$

This approach is limited but can handle BC's (and show how to handle BC's to maintain stability)