

### Regular Perturbation Series

The solution of a regular perturbation problem is usually a smooth function of some power  $q$  of  $\varepsilon$  (e.g.  $\varepsilon$ ,  $\varepsilon^2$  or  $\varepsilon^{1/2}$ ) near  $\varepsilon = 0$ . In this case, we can expand the solution in a Taylor series in  $\varepsilon^q$ , which we call a *perturbation series*:

$$y(x, \varepsilon) = y_0(x) + \varepsilon^q y_1(x) + \varepsilon^{2q} y_2(x) \dots \quad (2.1)$$

The function  $y_0(x)$  (the solution of the  $\varepsilon = 0$  problem) and the unknown functions  $y_1(x)$ ,  $y_2(x)$ , ... are found by substituting (2.1) into the original problem (including any boundary conditions), separating the terms according to powers of  $\varepsilon^q$ , and satisfying the successive problems at each order in  $\varepsilon$ . The exponent  $q$  is usually 1, but when the choice of  $q$  is unclear, it can be made by substituting the proposed solution into the original equation and seeing what  $q$  is required to balance the perturbation terms (the higher powers of  $\varepsilon$ ) in the equation. Like a normal Taylor series, the perturbation series is unique; if a series in powers of  $\varepsilon$  is assumed for  $y(x, \varepsilon)$ , but  $y(x, \varepsilon)$  could actually be written as a series in powers of  $\varepsilon^2$ , then the coefficients of the odd powers of  $\varepsilon$  will all turn out to be identically zero. However, if  $y(x, \varepsilon)$  really were a series in powers of  $\varepsilon^{1/2}$ , then our assumed series in powers of  $\varepsilon$  would be inconsistent. That is, we would find some orders of  $\varepsilon$  gave unbalanced equations where a nonzero function must equal zero. This would indicate that  $q$  should be changed to remove these imbalances.

#### Example (R1)

Find perturbation series for the roots of  $p(r; \varepsilon) = r^2 + \varepsilon r - 1 = 0$ ,  $\varepsilon \ll 1$

Note that for this problem, the two roots can be found exactly for any  $\varepsilon$  using the quadratic formula:

$$r^\pm(\varepsilon) = \frac{-\varepsilon \pm (4 + \varepsilon^2)^{1/2}}{2} \quad (2.2)$$

so this is just a pedagogical exercise to get us started. We can see that the two roots are smooth functions of  $\varepsilon$  for small  $\varepsilon$  and we can anticipate from (2.2) that the perturbation series will be the Taylor series in  $\varepsilon$  that we could derive from (2.2), and that because  $r^\pm(\varepsilon)$  are analytic functions of the complex variable  $\varepsilon$  out to the nearest singularities (branch points at  $\varepsilon = \pm 2i$ ), these series will have a radius of convergence of 2, i. e. they will converge for  $|\varepsilon| < 2$ . Let's see if this really works!

Let us try a perturbation series of the form

$$r(\varepsilon) = r_0 + \varepsilon r_1 + \varepsilon^2 r_2 \dots$$

and substitute into the polynomial:

$$\begin{aligned} 0 &= r^2 + \varepsilon r - 1 = (r_0 + \varepsilon r_1 + \varepsilon^2 r_2 \dots)^2 + \varepsilon(r_0 + \varepsilon r_1 + \varepsilon^2 r_2 \dots) - 1 \\ &= (r_0^2 - 1) + \varepsilon(2r_0 r_1 + r_0) + \varepsilon^2(2r_0 r_2 + r_1^2 + r_1) \dots \end{aligned}$$

For this to hold for arbitrary small values of  $\varepsilon$ , coefficients of all powers of  $\varepsilon$  must vanish:

$$\varepsilon^0: \quad 0 = r_0^2 - 1 \quad \Rightarrow \quad r_0^\pm = \pm 1$$

$$\varepsilon^1: \quad 0 = 2r_0 r_1 + r_0 \quad \Rightarrow \quad r_1^\pm = -\frac{1}{2}$$

$$\varepsilon^2: \quad 0 = 2r_0 r_2 + r_1^2 + r_1 \quad \Rightarrow \quad r_2^\pm = -\frac{r_1^2 + r_1}{2r_0} = -\frac{-1/4}{\pm 2} = \pm \frac{1}{8}$$

so we obtain the perturbation series for the two roots:

$$r^+(\varepsilon) = 1 - \frac{1}{2}\varepsilon + \frac{1}{8}\varepsilon^2 \dots$$

$$r^-(\varepsilon) = -1 - \frac{1}{2}\varepsilon - \frac{1}{8}\varepsilon^2 \dots$$

The series for  $r^+(\varepsilon)$  is indeed the same as the Taylor series of the exact solution around 0:

$$\frac{-\varepsilon + (4 + \varepsilon^2)^{1/2}}{2} = \frac{-\varepsilon}{2} + \left(1 + \frac{\varepsilon^2}{4}\right)^{1/2} = \frac{-\varepsilon}{2} + 1 + \frac{\varepsilon^2}{8} \dots$$

and similarly for  $r^-(\varepsilon)$ .

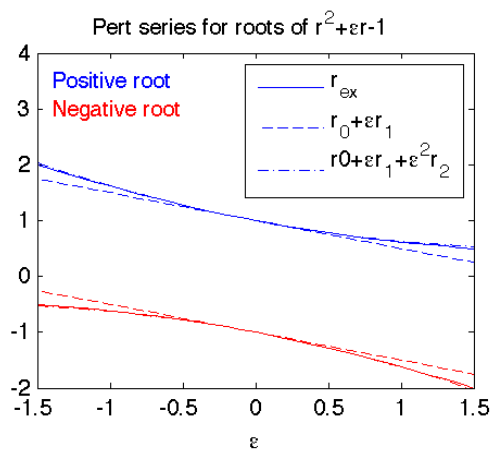


Fig. 2.1: Convergence of perturbation series for (R1)

In this case, the perturbation series are quite accurate even for fairly large values of  $\varepsilon$  (Fig. 2.1). If only the  $O(\varepsilon)$  term is kept in the series for either root, the  $O(\varepsilon^2)$  term gives an error estimate of roughly  $\varepsilon^2/8$ . If the  $O(\varepsilon^2)$  term is kept in the perturbation series, it remains highly accurate out to  $|\varepsilon| \approx 1.5$ .

### A trickier root-finding example

We wish to find perturbation series for the roots of the cubic equation

$$p(r; \varepsilon) = r^3 + r^2 - \varepsilon = 0, \quad \varepsilon \ll 1. \quad (2.3)$$

Its real roots (but not any complex roots) can be visualized by plotting  $p(r; \varepsilon)$  as in Fig. 2.2; the roots are where it crosses the  $r$ -axis. From the graph, we can see that the unperturbed problem has a root  $r_0^I = -1$  and a double root  $r_0^{II,III} = 0$  (we can tell it is a double root because  $dp/dr$  is also zero at  $r = 0$ ). The double root of the unperturbed problem is what makes the perturbation series more tricky. Graphically, we can see that for a small positive  $\varepsilon$ ,  $p(r; \varepsilon)$  locally looks like an upward facing parabola around its minimum value  $-\varepsilon$  at  $r = 0$ . Thus we can anticipate that the two roots  $r^{II,III}(\varepsilon)$  will be roughly  $\pm O(\varepsilon^{1/2})$ , rather than being linear functions of  $\varepsilon$ . On the other hand, the simple root  $r^I(\varepsilon)$  will be  $-1 + O(\varepsilon)$ . Let's explore how this emerges from the perturbation series approach!

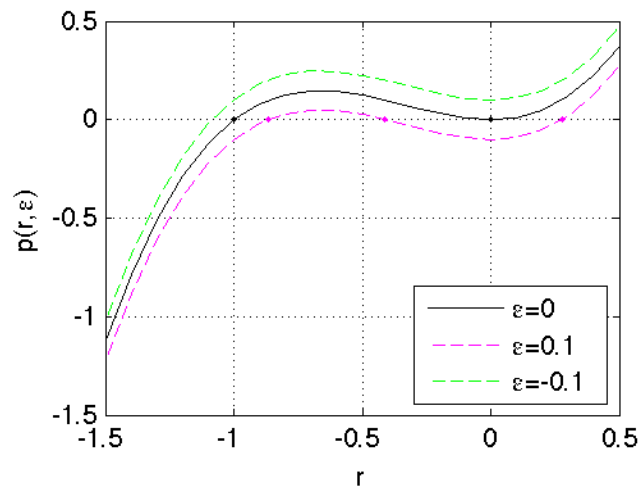


Fig. 2.2: Graph of  $p(r; \varepsilon)$ ; roots marked as dots for  $\varepsilon = 0$  and  $0.1$ .